

The Crystal and Molecular Structure of *N,N,N',N'*-Tetramethyl-*p*-diaminobenzene Iodide (Wurster's Blue Iodide)

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Crystals of Wurster's blue iodide have orthorhombic symmetry, with $a=5.919$, $b=9.855$, $c=9.901$ Å, space group $Pn\bar{m}$, $Z=2$. Intensities were collected with an automatic Nonius diffractometer. Atomic coordinates were obtained from Fourier and difference Fourier maps. Anisotropic least-squares refinement gave $R=2.27\%$ for 821 independent reflexions with reliable intensity. The *N,N,N',N'*-tetramethyl-*p*-diaminobenzene (TMPD) molecule lies at a special position with symmetry $2/m$. The nitrogen atom is at a distance of 0.095 Å from the plane through its surrounding carbon atoms. The benzene ring is planar. The *p*-diaminobenzene group has a quinoid character as may be seen from the bond lengths C(1)–C(1) 1.361, e.s.d. 0.008 Å; C(1)–C(2) 1.422, e.s.d. 0.005 Å; C(2)–N 1.344, e.s.d. 0.006 Å. The hydrogen atoms could be clearly identified in a difference Fourier synthesis. After least-squares refinement the average C–H distance was 0.965 Å (s.d. calculated by the least-squares program 0.04 Å for the individual values).

Introduction

Wurster's blue iodide has been investigated as part of a program to study compounds containing *N,N,N',N'*-tetramethyl-*p*-diaminobenzene (TMPD) molecules. Among these compounds there are diamagnetic and paramagnetic species and it seemed interesting to find out in which respect the structures of the two types of compound differ. Some compounds, e.g. Wurster's blue perchlorate, undergo a phase transition and are paramagnetic at high and diamagnetic at low temperatures. In Wurster's blue perchlorate the phase transition occurs at -87°C . Wurster's blue iodide is paramagnetic from room temperature till liquid nitrogen temperature, and no phase transition has been observed in this temperature interval.

First Wurster's blue perchlorate was studied at room temperature by one of the authors (J.L.d.B.). Bond lengths and angles obtained for the TMPD group, however, possibly suffered from systematic errors as it appeared that the perchlorate ions are statistically distributed in the crystal. To avoid this difficulty, Wurster's blue iodide was chosen for further research. The iodide was preferred above the corresponding chloride and bromide, as the latter compounds are unstable. At the end of the investigation of the iodide we noticed that the systematic errors in the TMPD unit of the perchlorate had been kept small, as the TMPD units in the two compounds show good agreement.

Study of Wurster's blue perchlorate will be continued at liquid nitrogen temperature. This work is severely hampered, however, as the crystals are twinned below the transition point (Thomas, Keller & McConnell, 1963).

Experimental

Crystal data

Wurster's blue iodide, $\text{C}_{10}\text{H}_{16}\text{N}_2\text{I}$, F.W. 291.157. M.P. $150\text{--}153^\circ\text{C}$. Orthorhombic, $Pn\bar{m}$, $a=5.919(3)$, $b=9.855(4)$, $c=9.901(4)$ Å. $V=577.5(5)$ Å³, $Z=2$, $D_m=1.68(4)$ g.cm⁻³ (flotation method), $D_c=1.674$ g.cm⁻³. Linear absorption coefficients: $\mu(\text{Cu } K\alpha)=219.1$ cm⁻¹, $\mu(\text{Mo } K\alpha)=27.66$ cm⁻¹. Crystals of ap-

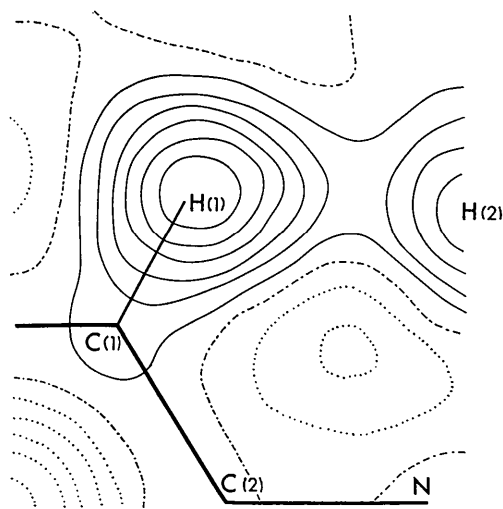


Fig. 1. Difference Fourier synthesis of low order reflexions in the plane of the benzene ring (approximated by the plane $-0.55431 X + 0.83231 Y + 0 Z = 0$ Å, with $X=xa$, $Y=yb$ and $Z=zc$). Contours at intervals 0.07 e.Å⁻³. The zero line is chain-dotted. The numbering of the atoms refers to Fig. 4, their positions are those of Table 1.

proximate dimensions of 0.09 mm perpendicular to the rotation axis were used for the determination of the cell constants, which were obtained from zero layer Weissenberg photographs, superimposed with NaCl reflexion spots for calibration purposes. The standard deviations given above are twice the standard deviations as calculated by the least-squares program used for the adjustment of a , b and c to the $\sin^2 \theta / \lambda^2$ values of 182 reflexions ($\lambda(\text{Cu } K\alpha_2) = 1.54434$, $\lambda(\text{Cu } K\alpha_1) = 1.54050$ Å, $a(\text{NaCl}) = 5.64006$ Å).

The systematic absences on the X-ray photographs were characteristic of both the space groups $Pnmm$ and $Pnm2$. Space group $Pnmm$, having the general positions $\pm(x, y, z)$; $-x, -y, z$; $\frac{1}{2} + x, \frac{1}{2} - y, \frac{1}{2} - z$; $\frac{1}{2} - x, \frac{1}{2} + y, \frac{1}{2} - z$

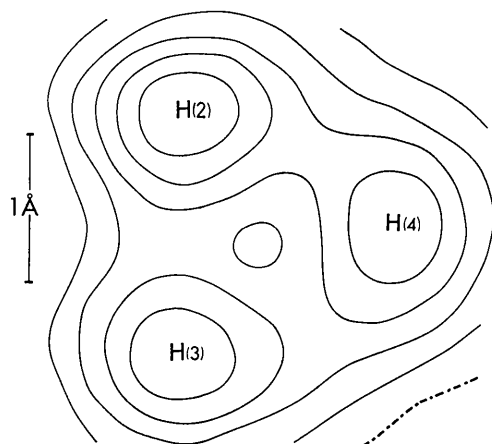


Fig. 2. Difference Fourier synthesis of low order reflexions in the plane of the methyl hydrogen atoms (approximated by the plane $0.38289 X + 0.45994 Y + 0.80115 Z = 3.421$ Å, with $X = xa$, $Y = yb$ and $Z = zc$). Contours at intervals 0.07 e.Å^{-3} . The zero line is chain-dotted. The numbering of the atoms refers to Fig. 4.

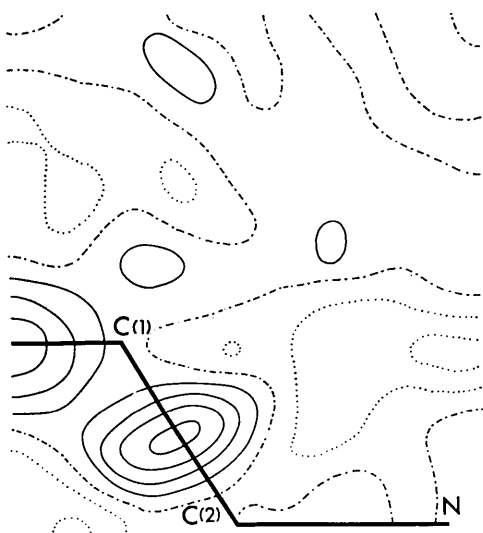


Fig. 3. Final difference Fourier synthesis (real part of ΔF only) in the plane of the benzene ring. Details are as in legend of Fig. 1.

was adopted for the structure determination. In this space group, both the iodine atoms and the TMPD groups lie at special positions with symmetry $2/m$. The results of the structure determination subsequently confirmed the choice of the centrosymmetric space group.

Preparation of the crystals

The substance was prepared by mixing, at room temperature, solutions of iodine (0.16 g.at.l^{-1}) and of TMPD (0.32 mol.l^{-1}) in acetonitrile in a 2:1 ratio. By slow evaporation of the acetonitrile in about 2 days, purple-black crystals, prismatic along a , could be obtained. A great number of the larger crystals appeared to contain mother liquor in channels parallel to the a axis. Several crystals were cut perpendicular to this axis to select a good specimen. For the intensity measurements, part of a crystal without channels in the section perpendicular to a , was used. The approximate dimensions of this crystal fragment are $0.39 \times 0.42 \times 0.44 \text{ mm}^3$.

Structure factors

The crystal was mounted on the tip of a glass fibre and aligned along a . The intensity data were collected with an automatic Nonius diffractometer, using the moving counter-moving crystal method and Zr-filtered Mo radiation. Deviations from linearity of the (scintillation) counting equipment were kept below 1% by the use of attenuation filters. Every 90 minutes a reference reflexion was measured accurately. Two equivalent sets of independent reflexions, hkl and $h\bar{k}l$, with $\sin \theta / \lambda < 0.8171 \text{ Å}^{-1}$, were subsequently measured. Corrections were made for intensity changes in the primary beam (by use of the reference intensities), for the Lorentz and polarization effects and for absorption. The latter corrections were calculated according to the Busing & Levy (1957) scheme; $6 \times 6 \times 6$ volume fragments and 11 accurately located bounding planes were taken into account. In the part of reciprocal space explored, 1381 independent reflexions were measured. From both sets of measurements were omitted (1) reflexions with a negative net intensity (background > integrated intensity) and (2) reflexions with a weight w smaller than a seventh of the average weight of the reflexions with positive net intensity. The weights were calculated from the relation $w = \sigma^{-2}(F)$, $\sigma(F)$ being the standard deviation in F due to counting statistics.

Set 1 (hkl) contained 917 reliable reflexions, set 2 ($h\bar{k}l$) 903. The two sets had 823 reliable reflexions in common; 174 reflexions were above the threshold weight in only one of the two sets. From a comparison of the 823 reliable reflexions common to both sets, the scaling factor K calculated according to $K = \Sigma |F(\text{set 1})| / \Sigma |F(\text{set 2})|$ amounted to 1.002 and the discrepancy $\Sigma |\Delta F| / \Sigma |F|$ to 2.4%.

The structure determination was carried out with the 823 common reflexions. As $F_o(hkl)$, the average F value of hkl and $h\bar{k}l$ was taken.

Determination of the structure

Both the TMPD group and the iodine atom are situated at twofold positions with symmetry $2/m$. In order to facilitate comparison with Wurster's blue perchlorate (Thomas *et al.*, 1963), the iodine atoms were located at $(0, \frac{1}{2}, 0)$ and $(\frac{1}{2}, 0, \frac{1}{2})$. Apart from possible anisotropic vibrations, the iodine atoms contribute only to reflexions hkl with $h+k+l=\text{even}$. A Fourier synthesis based on this set of reflexions, with signs deduced from the iodine contributions alone, was calculated. The synthesis showed the TMPD molecule around $(0, 0, 0)$ and the related one at $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$. For each molecule two images were found as a result of the pseudo-symmetry introduced in the Fourier map owing to the fact that only reflexions hkl with $h+k+l=\text{even}$ were considered. One image was arbitrarily regarded as the spurious one and the positions of the one independent nitrogen and the three independent carbon atoms could be determined from the map.

The parameters of I, C and N were refined with the least-squares method in which the scattering factors were taken from Moore (1963) and anisotropic thermal parameters were considered. The calculations were done on a TR4 computer with a program according to

Cruickshank's (1961*b*) scheme. The weighting scheme adopted was $w=710/(710+F_o^2)$ and R decreased to 3.4%. To find the hydrogen atoms, a difference Fourier synthesis of the reflexions with $\sin \theta/\lambda < 0.4 \text{ \AA}^{-1}$ (H not included in F_o) was calculated in the plane of the benzene ring and in the plane through the hydrogen atoms of the methyl group; the latter plane was calculated with the assumption that the C-H distances are 1.08 \AA and that the angles at the methyl carbon atom are tetrahedral. Figs. 1 and 2 clearly show the four independent H atoms. The hydrogen atoms were included in the least-squares refinement, with isotropic parameters. For these atoms the scattering factor as given by Stewart, Davidson & Simpson (1965) was used. In the latter stage of the refinement the anomalous part of the iodine scattering factor, $\Delta f' = -0.6$, $\Delta f'' = 2.2$ (*International Tables for X-ray Crystallography*, 1962), was taken into account and the weighting scheme $w=525/(525+F_o^2)$ was adopted. The reflexions 002 and 121, which clearly suffered from extinction, were not considered during the refinement. In the last cycle the changes in the parameters of I, N and C atoms were smaller than 0.05 times the standard deviations as calculated by the least-squares program, those of the hydrogen atoms were smaller than 0.1 of

Table 1. *Final coordinates and isotropic temperature factors*

x, y, z are obtained from the least-squares refinement; x', y', z' are corrected for libration. For numbering of atoms see Fig. 4.

I	x 0	y $\frac{1}{2}$	z 0	B	x'	y'	z'
C(1)	0.09484 (44)	0.03907 (23)	0.12225 (21)		0.09501	0.03910	0.12297
C(2)	0.19763 (51)	0.08290 (29)	0		0.19801	0.08298	0
C(3)	0.47241 (39)	0.22005 (30)	0.12391 (29)		0.47306	0.22028	0.12454
N	0.38209 (46)	0.16244 (27)	0		0.38276	0.16259	0
H(1)	0.15741 (432)	0.06443 (265)	0.20680 (261)	3.5 (5) \AA^2			
H(2)	0.48628 (451)	0.15065 (468)	0.19696 (431)	6.5 (10)			
H(3)	0.62502 (616)	0.25274 (333)	0.10945 (341)	6.6 (8)			
H(4)	0.37493 (604)	0.28879 (332)	0.15338 (354)	6.4 (8)			

Table 2. *Thermal parameters of the temperature factor*

$$\exp [-2\pi^2 (h^2 a^2 U_{11} + k^2 b^2 U_{22} + l^2 c^2 U_{33} + 2ha*kb*U_{12} + 2kb*lc*U_{23} + 2ha*lc*U_{13})]$$

I	U_{11}	U_{22}	U_{33}	$2U_{12}$	$2U_{23}$	$2U_{13}$
I	0.0540 (2)	0.0590 (2)	0.0515 (2)	0.0077 (3)	0	0
C(1)	0.0506 (11)	0.0449 (9)	0.0316 (8)	0.0015 (19)	-0.0017 (15)	-0.0054 (18)
C(2)	0.0399 (13)	0.0359 (12)	0.0350 (12)	0.0128 (22)	0	0
C(3)	0.0484 (14)	0.0626 (14)	0.0522 (12)	-0.0165 (21)	-0.0070 (23)	-0.0165 (20)
N	0.0393 (12)	0.0424 (12)	0.0421 (12)	0.0042 (21)	0	0

Table 3. $T(\text{\AA}^2)$ and $\omega(\text{rad}^2)$ tensors and the corresponding 'calculated' thermal parameters

p, q and r are the main axes for T ; p', q' and r' those for ω . r and r' coincide with c ; p and p' are obtained by rotating the x axis about c over angles of 40.03° and 40.37° respectively in the direction of y .

	T_p	T_q	T_r	$\omega_{p'}$	$\omega_{q'}$	$\omega_{r'}$
	0.0443	0.0276	0.0304	0.0121	0.0017	0.0015
	U_{11}	U_{22}	U_{33}	$2U_{12}$	$2U_{23}$	$2U_{13}$
C(1)	0.0465	0.0463	0.0313	0.0008	-0.0002	-0.0035
C(2)	0.0384	0.0366	0.0341	0.0136	0	0
C(3)	0.0537	0.0581	0.0520	-0.0175	-0.0060	-0.0143
N	0.0413	0.0423	0.0441	0.0054	0	0

their standard deviations. The value of $\langle w(\Delta F)^2 \rangle$ appeared to be approximately independent of $|F|$. The final R index is 2.27% for the 821 reflexions.

The final parameters with the standard deviations as calculated by the least-squares program are listed in Tables 1 and 2. The final difference map in the plane

of the benzene ring is shown in Fig. 3. It may be seen that there are maxima between the bonded atoms of the benzene ring. These maxima are about 5 times the standard deviation, $0.06 \text{ e.}\text{\AA}^{-3}$, obtained by Cruickshank's (1949) method and may be ascribed to bonding effects. Apart from these maxima, the difference map

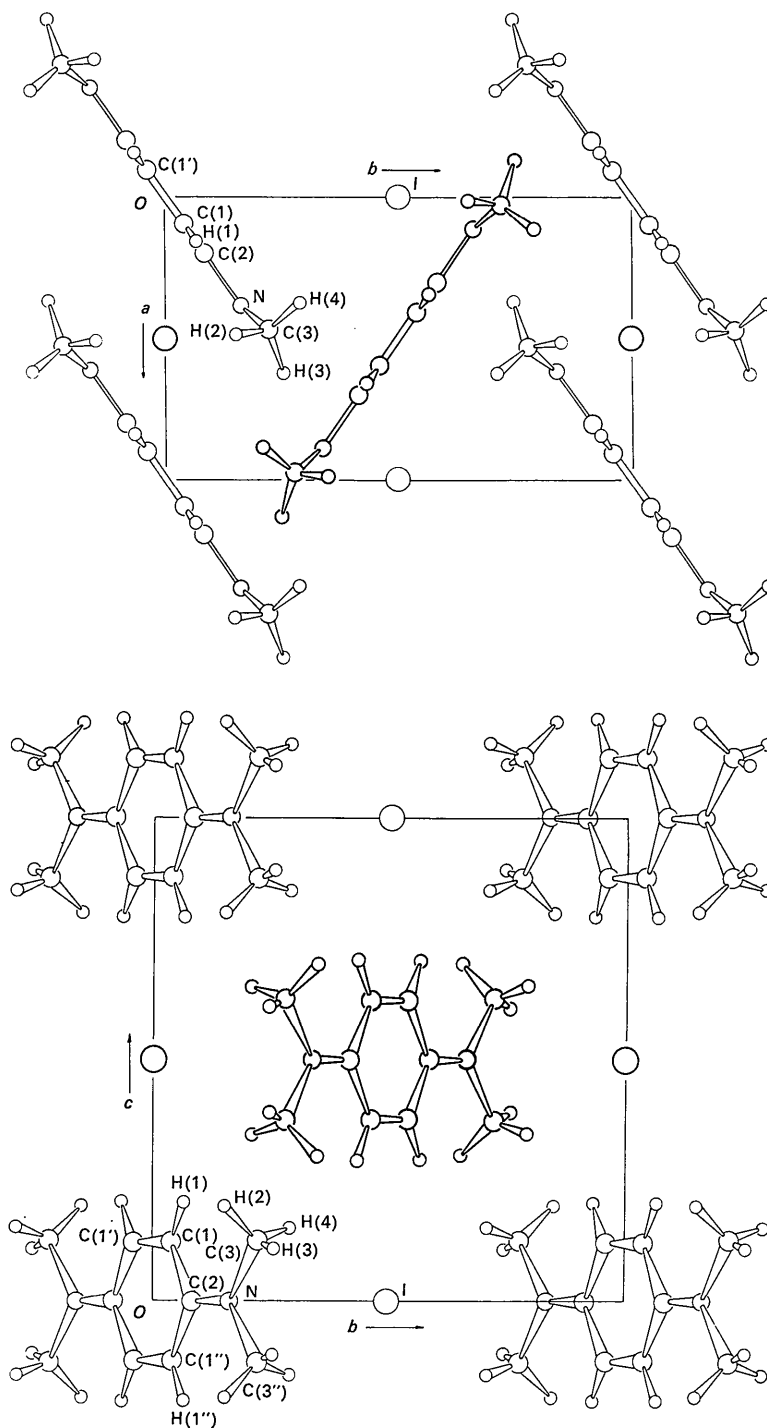


Fig. 4. Projections of the structure along [001] and [100].

shows some maxima (four times the e.s.d.) at distances of 0.8 Å from the iodine atom. The map shows negative values (up to 5 times the e.s.d.) in a region related to the position of the molecule by a translation $\frac{1}{2}, \frac{1}{2}, \frac{1}{2}$. No

significant features are further found in the difference map.

It is noteworthy that the final positions of the hydrogen atoms are approximately 0.96 Å from the car-

Table 4. Observed and calculated structure factors

In order to obtain *F* on the absolute scale, the values in the Table must be divided by 7.0.

H K L FO FC AC BC												H K L FO FC AC BC												H K L FO FC AC BC												H K L FO FC AC BC												H K L FO FC AC BC											
0	0	0	212221	2731272	42	1	2	3	2227	221-221	0	2	0	10	161	154	153	14	2	11	7	78	79	-78	-5	3	10	7	49	49	48	6	4	13	3	48	48	-45	-3	6	4	128	127	126	10														
0	0	0	475	472	472	10	1	2	3	42	48	-42	0	2	0	12	33	38	37	5	2	12	0	39	41	-40	-2	4	13	5	49	47	41	45	4	13	5	49	47	-41	-5	6	4	69	68	68	8												
0	0	0	668	662	662	10	1	2	4	133	117	-117	0	2	0	14	33	38	37	5	2	12	0	39	41	-40	-2	4	13	5	49	47	41	45	4	13	5	49	47	-41	-5	6	4	69	68	68	8												
0	0	0	898	892	892	10	1	2	5	383	392	391	2	1	0	133	134	134	0	2	12	0	60	58	58	7	3	11	2	93	93	-93	-8	5	0	1292	287	286	20	6	4	10	1	44	44	44	4												
0	0	0	1027	1021	1021	10	1	2	6	1073	101	-101	0	2	1	1	708	730	-729	-38	2	12	4	70	68	67	6	3	11	3	32	27	27	1	5	0	3433	401	400	19	6	5	0	57	58	58	1												
0	0	0	1275	1277	1276	10	1	2	7	501	505	502	24	2	1	2	45	46	45	0	2	12	6	42	39	38	5	11	4	102	99	-99	-7	5	0	1333	337	337	15	6	5	1	131	131	-111	-11													
0	0	0	1475	1475	1475	10	1	2	8	91	94	-94	0	2	1	3	787	790	-790	-35	2	13	1	61	63	-62	-5	3	11	6	69	67	-67	-5	5	0	1757	156	156	12	6	5	2	29	23	23	1												
0	0	0	1675	1673	1673	10	1	2	9	335	336	336	18	2	1	4	68	57	57	0	2	13	3	57	55	-55	-5	3	12	1	71	67	66	6	5	0	88	88	88	9	6	5	122	124	-124	-10													
0	0	0	1875	1874	1874	10	1	2	10	52	49	-49	0	2	1	5	597	599	-599	-30	2	13	4	30	9	9	0	3	12	3	40	42	42	5	5	0	1105	102	102	6	6	5	5	107	105	-105	-8												
0	0	0	2075	2075	2075	10	1	2	11	194	196	193	2	2	1	6	72	43	-43	0	2	13	6	52	47	-47	-4	4	12	5	36	33	32	5	5	0	137	74	73	4	6	5	7	77	75	-75	-7												
0	0	0	2275	2275	2275	10	1	2	12	31	31	31	0	2	1	7	315	335	-334	-23	2	13	7	40	38	-38	-3	3	12	6	27	27	27	0	0	1	302	308	-307	-20	6	5	9	47	49	-49	-5												
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0	0	0	2675	2675	2675	10	1	2	14	115	115	115	0	2	1	9	115	156	-156	-11	2	14	3	442	447	446	33	3	12	8	49	51	-51	-4	6	1	3	63	61	-61	-6	6	0	88	89	88	10												
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0	0	0	3875	3875	3875	10	1	2	20	589	589	589	0	2	1	15	224	180	180	18	2	17	9	628	180	180	18	3	14	10	49	29	29	0	0	2	4	293	296	295	23	6	7	1	107	107	-107	-8											
0	0	0	4075	4075	4075	10	1	2	21	668	668	668	0	2	1	16	241	196	196	19	2	18	10	707	196	196	19	3	14	11	58	38	38	0	0	3	257	256	256	18	6	7	3	85	85	-85	-7												
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0	0	0	5475	5475	5475	10	1	2	28	1221	1221	1221	0	2	1	23	361	308	308	28	2	22	17	900	308	308	28	3	14	18	132	112	112	-12	2	10	174	174	174	12	6	7	10	10	10	10	10	10	10										
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0	0	0	6275	6275	6275	10	1	2	32	1537	1537	1537	0	2	1	27	429	376	376	32	2	24	21	976	376	376	32	3	14	22	188	168	168	-16	2	12	224	224	224	12	6	7	14	14	14	14	14	14	14										
0	0	0	6475	6475	6475	10	1	2	33	1616	1616	1616	0	2	1	28	446	393	393	33	2	24	22	995	393	393	33	3	14	23	206	186	186	-16	2	13	238	238	238	12	6	7	15	15	15	15	15	15	15										
0	0	0	6675	6675	6675	10	1	2	34	1695	1695	1695	0	2	1	29	463	410	410	34	2	24	23	1014	410	410	34	3	14	24	224	204	204	-16	2	13	252	252	252	12	6	7	16	16	16	16	16	16	16										
0	0	0	6875	6875	6875	10	1	2	35	1774	1774	1774	0	2	1	30	480	427	427	35	2	25	24	1033	427	427	35	3	14	25	242	222	222	-16	2	14	266	266	266	12	6	7	17	17	17	17	17	17	17										
0	0	0	7075	7075	7075	10	1	2	36	1853	1853																																																

bon atoms to which they are attached, whereas visual inspection of the difference map of the low angle reflexions (see contour lines of H(1) in Fig. 1) suggests that the hydrogen atom H(1) lies close to its expected position at 1.08 Å from C(1). The shifts of the hydrogen atoms towards the carbon atoms during the least-squares refinement were assumed to be due to bonding effects being more pronounced when high order reflexions are taken into account. Attempts to determine the positions of the hydrogen atoms from the difference map of the low order reflexions failed, however, as their maxima are very flat. Least-squares refinement of the reflexions with $\sin \theta/\lambda < 0.4 \text{ \AA}^{-1}$ with fixed parameters for I, N and C, again gave C-H distances of about 0.96 Å. The final parameters of the least-squares refinement with the 821 reflexions were therefore taken as the best parameters for the hydrogen atoms.

To check whether the TMPD molecule may be regarded as a rigid body, the thermal parameters of the C and N atoms were analysed according to Cruickshank's (1956) method. The values of T and ω are given in Table 3. Comparison of the 'experimental' and 'calculated' thermal parameters (Tables 2 and 3 respective-

ly) shows that the rigid body assumption is not strictly true for our molecule. Therefore only approximate corrections for libration effects (Cruickshank, 1961*a*) could be obtained. The corrected coordinates are listed in Table 1. The further discussion is based on the uncorrected values. Use of the corrected values would have increased the bond lengths by 0.002 to 0.007 Å.

At the end of the structure refinement, the 174 reflexions which were above the threshold weight in only one of the two sets of measurements, were considered. As F_0 , the average F value of hkl and $h\bar{k}l$ was taken, negative net intensities being considered as zero. The atomic shifts caused by inclusion of the 174 reflexions in the least-squares refinement were largest for C(1) and such that the length of the central bond of the benzene ring increases with 0.007 Å. For the remaining atoms the shifts are very small. From the standard deviations calculated by the least-squares program it appeared that the parameters of this refinement are less accurate than the final parameters in Tables 1 and 2, they were therefore not considered further.

The observed and calculated structure factors, as based on the parameters of Tables 1 (x, y, z) and 2, are listed in Table 4.

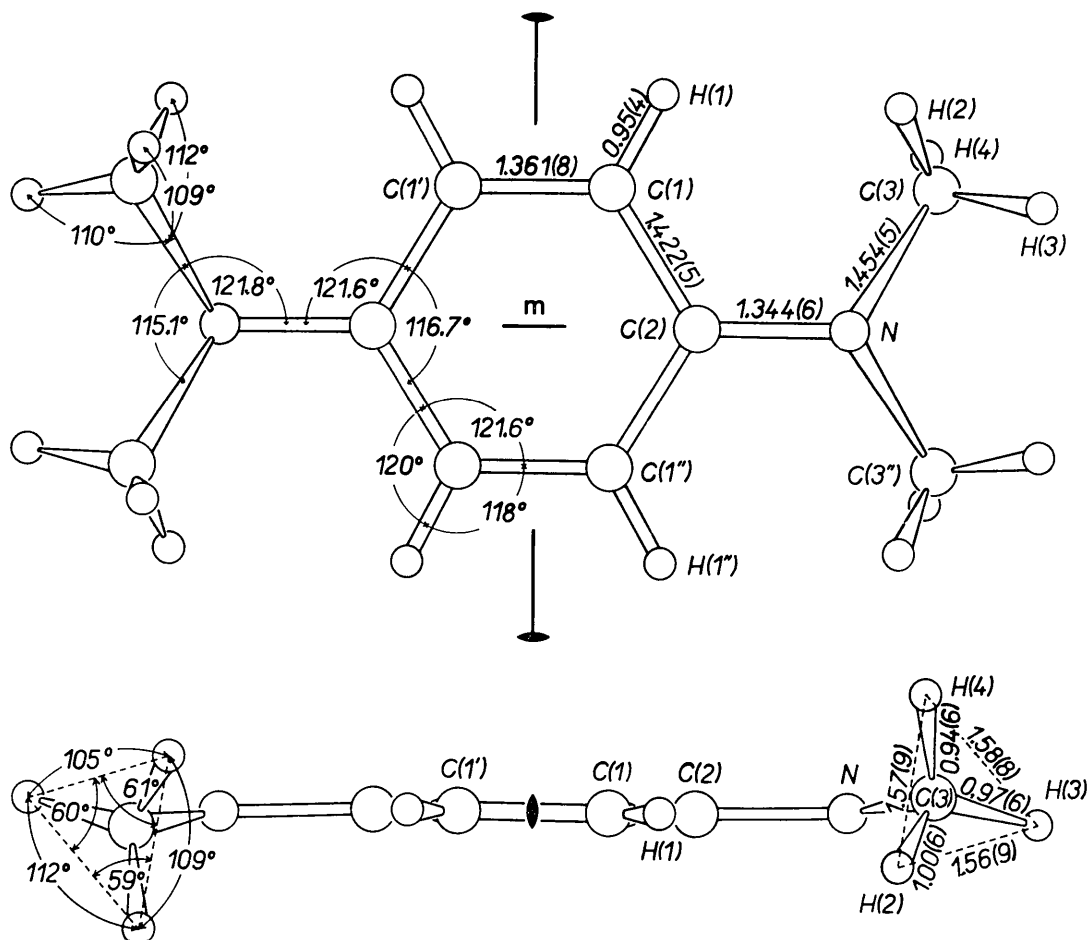


Fig. 5. Distances and angles in the TMPD molecule. The numbering of the atoms refers to Fig. 4.

Discussion of the structure

Projections of the structure along the [001] and [100] axes are shown in Fig. 4. The shortest intermolecular distances are given in Table 5. It appears that none of the intermolecular contacts is shorter than expected from the van der Waals radii (2.15, 1.7, 1.5 and 1.0 Å for I, C, N and H respectively (Nyburg, 1961)).

Table 5. *Intermolecular distances*

Only distances exceeding the sum of the van der Waals radii by less than 0.5 Å are given. For numbering of atoms see Fig. 4; C(3), *r* means C(3) at $1-x, -y, z$; C(1''), *n* means C(1'') at $\frac{1}{2}+x, \frac{1}{2}-y, \frac{1}{2}+z$; etc.

I(0½0)···C(2)	4.274 Å	C(1)···C(3), <i>r</i>	3.617 Å
I(0½0)···N	4.023	C(1)···H(2), <i>r</i>	3.192
I(0½0)···C(3)	4.115	C(2)···C(3), <i>r</i>	3.773
I(0½0)···H(4)	3.400	C(2)···N, <i>r</i>	3.469
I(¼0½)···H(1)	3.598	N···N, <i>r</i>	3.493
I(¼0½)···C(3)	4.312	C(3)···C(1''), <i>n</i>	3.532
I(¼0½)···H(2)	3.349	C(3)···H(1''), <i>n</i>	2.919

The TMPD molecule with bond lengths and angles is shown in Fig. 5; the standard deviations in the bond lengths are given in parentheses. The standard deviations in angles involving H atoms range from 2.6° to 5.0°, those of the remaining angles are 0.4° on the average. In calculating the standard deviations in the bond lengths and angles, the s.d. in the coordinates (Table 1), which are lower estimates, were multiplied by 1.6. The benzene ring is planar within experimental error, the atoms C(1) and C(2) lie only at distances of 0.005 and 0.002 Å respectively from the best plane through the ring. The C(2)–N bond and the N–C(3)–C(3'') plane are bent away from the benzene plane by 0.83° and 11.88° respectively.

Although, strictly speaking, the molecule is not a rigid body, a good impression of its thermal movement could be obtained. From the following data it appears that the translation and rotation are largest along axes which almost coincide with the direction in which the molecule is elongated most. The plane of the benzene ring makes an angle of 34.8° with the *a* axis, whereas the angle between the plane N–C(3)–C(3'') and the *a* axis is 46.7°. The main axes *p* and *p'* for which the components of *T* and ω are largest make angles of 40.0° and 40.4° respectively with the *a* axis. T_p is about 1.5 times the remaining *T* components, which are nearly equal, $\omega_{p'}$ exceeds the average values of the remaining ω components by as much as a factor 7.5.

The C–C bond lengths in the TMPD unit in Wurster's blue iodide, C(1')–C(1) 1.361, C(1)–C(2) 1.422 Å, may be compared with the lengths obtained by theoretical calculations for the neutral TMPD molecule and the positive TMPD⁺ ion, 1.388 and 1.402 Å respectively for TMPD⁰ and 1.370 and 1.420 Å respectively for TMPD⁺. The π -bond orders *p* were calculated by the PPP method (Monkhorst & Kommandeur, 1967); to obtain the bond lengths *d* the equation $d = 1.51 - 0.18p$ was used. There appears to be good agreement be-

tween the experimental values and those calculated for the TMPD⁺ ion. As expected, TMPD iodide is thus TMPD⁺I⁻.

In Fig. 6 the TMPD⁺ ion is compared with the TMPD unit in TMPD chloranil, as recently determined by some of the authors (de Boer & Vos, 1968). The X-ray work on the latter compound indicates that the majority of the molecules is neutral although the presence of some 20% of doubly charged units, as required by spectral measurements (Pott, 1966), could not be excluded. The figure clearly shows that the TMPD⁺ ion in TMPD iodide has a more pronounced quinoid character than the 'neutral' TMPD unit in TMPD chloranil. In both TMPD units the bonds N–CH₃ are considerably longer than the bonds N–C₆H₄ which have a partial double bond character. The difference in length between the N–C₆H₄ bonds of the two molecules is in qualitative agreement with the calculated π bond orders, 0.417 for TMPD⁰ and 0.589 for TMPD⁺. The 'single' N–CH₃ bonds are, on the average, somewhat shorter than the value 1.47 Å quoted in *International Tables for X-ray Crystallography* (1962) for a single N–C(*sp*³) bond.

Consideration of the bond angles in the TMPD units shows that the angles have a tendency to increase with increasing double bond character of the adjacent bonds. A similar behaviour has been observed in other compounds (Gillespie & Nyholm, 1957).

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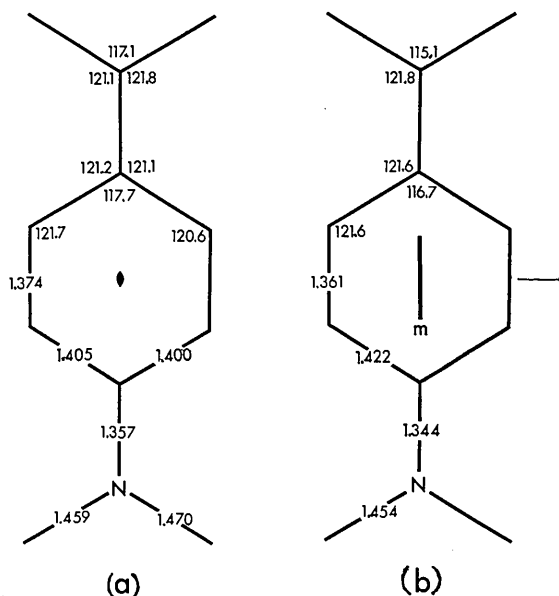


Fig. 6. Comparison of TMPD in TMPD chloranil (a) with TMPD⁺ in TMPD iodide (b). TMPD in TMPD chloranil is planar because of its special position in the unit cell.

of the cell constants and Drs R. Olthof-Hazekamp for writing and modifying many of the programs used. We are indebted to Dr K. Toman, Dr B. Sedlaček and Prof. E. H. Wiebenga for their interest in this work. We thank the staff of the Computing Centre of the University, Groningen, for their cordial cooperation.

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Hydrogen Bond Studies. XVII.* The Crystal Structure of Dimethylammonium Chloride

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The crystal structure of dimethylammonium chloride, $(\text{CH}_3)_2\text{NH}_2\text{Cl}$, has been determined from three-dimensional single-crystal X-ray data obtained at room temperature. The crystals are orthorhombic, space group *Ibam*, with eight formula units in the cell of dimensions $a = 7.301$, $b = 14.515$, $c = 9.926$ Å. The structure contains infinite planar chains of chlorine and nitrogen atoms linked together by hydrogen bonds of lengths 3.12 and 3.09 Å. The dimethylammonium and chloride ions form a distorted NaCl-type structure.

Introduction

In conjunction with the study of simple hydrogen-bonded structures in progress at this Institute an investigation of the changes in the N-H...Cl bond lengths accompanying replacement of hydrogen in ammonium by methyl groups is of considerable interest. The crystal structure of monomethylammonium chloride was studied by X-ray diffraction technique by Hughes & Lipscomb (1946).

In the present investigation the crystal structure of dimethylammonium chloride has been determined from single-crystal X-ray diffraction data obtained at 20°C.

Experimental

The sample used was commercial dimethylammonium chloride. The salt was recrystallized three times from chloroform, and dried in a vacuum desiccator as described by Whitlow & Felsing (1944). The crystals had to be handled in a dry-box as they were highly hygroscopic. In the first attempts powder samples were sealed in glass capillaries and single crystals were grown by a

zone-melting technique, the growth of the crystals being followed with a polarizing microscope. The melting point was found to agree with that earlier published, 167°C (Hunter & Byrkit, 1932). Most of the crystals thus obtained had their crystallographic axes orientated at an angle of 45° to the capillary axis. It proved very difficult to obtain a crystal with suitable orientation. Weissenberg photographs demonstrated that the crystals were probably twinned; furthermore the X-ray reflexions became unusually diffuse on passing from lower to higher layers. Owing to these difficulties crystals grown from solution were tried instead. Plate-like crystals were obtained with chloroform alone as a solvent, but from a solution with equal amounts of chloroform and ethyl alcohol needle-like crystals with a nearly circular cross-section were obtained. One of these with a diameter of 0.15 mm and a length of approximately 2 mm was chosen and sealed in a thin-walled glass capillary. An appropriate crystallographic axis was thus obtained parallel to the capillary axis. With this, the *a* axis, as rotation axis, equi-inclination Weissenberg photographs, layers $0 \leq h \leq 6$, were recorded with Ni-filtered Cu $K\alpha$ radiation. It was now clear that the previous crystal, grown by the zone-melting procedure, was indeed a twin

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